

## Differential operators and basics of the Ricci flow

We recall that the square of the covariant derivative is

$$\nabla^2: \Gamma(T_g^p M) \rightarrow \Gamma(T_g^{p+2} M)$$

w/

$$\begin{aligned}\nabla^2 \alpha(X, Y, z_1, \dots, z_p) &= \nabla_X (\nabla \alpha)(Y, z_1, \dots, z_p) \\ &= [\nabla_X (\nabla \alpha(Y)) - \nabla \alpha(\nabla_X Y)] \\ &\quad (z_1, \dots, z_p) \\ &= \nabla_X (\nabla_Y \alpha)(z_1, \dots, z_p) - \nabla_{\nabla_X Y} \alpha(z_1, \dots, z_p)\end{aligned}$$

so in short

$$\nabla_{X, Y}^2 \alpha = \nabla_X \nabla_Y \alpha - \nabla_{\nabla_X Y} \alpha.$$

For ex. in local coordinates

$$\begin{aligned}\nabla_i R_{jk} &= (\nabla R_c)(\partial_i, \partial_j, \partial_k) \\ &= (\nabla_{\partial_i} R_c)(\partial_j, \partial_k)\end{aligned}$$

$$= \frac{\partial}{\partial x^i} R_{jk} - \Gamma_{ij}^l R_{lk} - \Gamma_{ik}^l R_{jl}$$

Similarly

$$\begin{aligned} \nabla_i R_{jklm} &= \partial_i R_{jklm} - \Gamma_{ij}^p R_{pklm} - \Gamma_{ik}^p R_{jplm} \\ &\quad - \Gamma_{il}^p R_{jkpm} - \Gamma_{im}^p R_{jkpe} \end{aligned}$$

for a function  $f$ , one can show that

$$\begin{aligned} \nabla_i \nabla_j f &= (\nabla \nabla f)(\partial_i, \partial_j) \\ &= \frac{\partial^2 f}{\partial x^i \partial x^j} - \Gamma_{ij}^k \frac{\partial f}{\partial x^k} \end{aligned}$$

\* We denote the  $m$ -th order derivative of  $\alpha$

$$\text{by } \nabla^m \alpha = \underbrace{\nabla \cdots \nabla}_m \alpha$$

$m$ -times

and its components will be denoted by

$$\nabla_{j_1 \cdots j_m} \alpha_{i_1 \cdots i_p}$$

## Lie Derivative

For  $X, Y \in \Gamma(TM)$

$$\mathcal{L}_X Y = [X, Y] \text{ w/}$$

$$[X, Y] \in \Gamma(TM) \text{ s.t.}$$

$$[X, Y](f) = X(Y(f)) - Y(X(f)).$$

If  $\theta$  is a 1-form then

$$(\mathcal{L}_X \theta)(Y) = X(\theta(Y)) - \theta([X, Y])$$

Even though one doesn't need a metric to define the Lie derivative, it is related to the L-C connection via the formula

$$(\mathcal{L}_X A)(Y_1, \dots, Y_p, \theta_1, \dots, \theta_q)$$

$$= X(A(Y_1, \dots, Y_p, \theta_1, \dots, \theta_q))$$

$$- \sum_{1 \leq i \leq p} A(Y_1, \dots, [X, Y_i], \dots, Y_p, \theta_1, \dots, \theta_q)$$

$$- \sum_{1 \leq j \leq q} A(Y_1, \dots, Y_p, \theta_1, \dots, \mathcal{L}_X \theta_j, \dots, \theta_q)$$

In particular, when  $A=g$  then

$$(\mathcal{L}_X g)(y, z) = g(\nabla_y X, z) + g(y, \nabla_z X)$$

whose expression in coordinates give

$$(\mathcal{L}_X g)_{ij} = \nabla_i X_j + \nabla_j X_i.$$

If  $X = \nabla f$  for  $f \in C^\infty(M)$  then

$$\begin{aligned} (\mathcal{L}_X g)_{ij} &= (\mathcal{L}_{\nabla f} g)_{ij} = \nabla_i \nabla_j f + \nabla_j \nabla_i f \\ &= 2 \nabla_i \nabla_j f. \end{aligned}$$

### Ricci Identities

we have

$$(\nabla_i \nabla_j - \nabla_j \nabla_i) \alpha_{R_1, \dots, R_n} = - \sum_{l=1}^n R_{ijkl}^m \alpha_{R_1, \dots, R_{l-1}, m, R_{l+1}, \dots, R_n}$$

So for a 1-form  $\alpha$ ,

$$(\nabla_i \nabla_j - \nabla_j \nabla_i) \alpha_R = - R_{ijkl} \alpha_l$$

or for a 2-tensor  $\beta$

$$\nabla_i \nabla_j \beta_{Rl} - \nabla_j \nabla_i \beta_{Rl} = - R_{ijkm} \beta_{ml} - R_{ijem} \beta_{km}.$$

The divergence of a  $(p,0)$ -tensor is

$$\begin{aligned}(\operatorname{div} \alpha)_{i_1 \dots i_{p-1}} &= g^{jk} \nabla_j \alpha_{k i_1 \dots i_{p-1}} \\ &= \nabla_j \alpha^j_{i_1 \dots i_{p-1}}\end{aligned}$$

for 1-forms  $\alpha$ ,  $\operatorname{div} \alpha$  is a function,

$$\operatorname{div} \alpha = \nabla_i \alpha^i.$$

### Laplacian

Laplacian  $\Delta$  on functions is  $\operatorname{div}(\operatorname{grad})$

i.e.,

$$\begin{aligned}\Delta &= \operatorname{div} \nabla = g^{ij} \nabla_i \nabla_j \\ &= g^{ij} \left( \frac{\partial^2}{\partial x^i \partial x^i} - \Gamma_{ij}^k \frac{\partial}{\partial x^k} \right)\end{aligned}$$

For tensors, again

$$\begin{aligned}\Delta &= \operatorname{div}(\operatorname{grad}) = \operatorname{trace}_g \nabla^2 \\ &= g^{ij} \nabla_i \nabla_j\end{aligned}$$

Exercise:- Prove the Bochner formula for  $|\nabla f|^2$

i.e.,  $\forall f \in C^\infty(M)$

$$\Delta |\nabla f|^2 = 2|\nabla \nabla f|^2 + 2R_{ij} \nabla_i f \nabla_j f + 2 \nabla_i f \nabla_i (\Delta f).$$

Conclude that if  $Rc \geq 0$ ,  $\Delta f = 0$  and  $|\nabla f| = 1$

then  $\nabla f$  is parallel.

We also have the divergence theorem and integration by parts formula. Let  $M^n$  be a closed manifold and  $u, v \in C^\infty(M)$  and  $X \in \Gamma(TM)$ .

Then

$$\int_M \operatorname{div} X \operatorname{vol} = 0 \quad \text{and so}$$

$$\int_M \Delta u \operatorname{vol} = 0$$

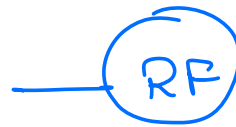
$$\int_M u \Delta v \operatorname{vol} = \int_M v \Delta u \operatorname{vol} \quad (\text{Integration by parts}).$$

## Ricci Flow

Let  $(M^n, g_0)$  be given.

Def<sup>n</sup> A Ricci flow on  $(M^n, g_0)$  is a family of metrics  $(g(t))_{t \in [0, \epsilon)}$  s.t.

$$\begin{aligned}\partial_t g(t) &= -2\text{Ric}(g(t)) \\ g(0) &= g_0.\end{aligned}$$



$\epsilon$  depends on  $M^n$  and  $g_0$ .

Examples :-

① If  $g_0$  is Ricci-flat, i.e.,  $\text{Ric} = 0 \Rightarrow g(t) = g_0$   $\forall t$  is a sol<sup>n</sup>. Note that  $t \in (-\infty, \infty)$  in this case. e.g. when  $M = \mathbb{R}^n$  or flat torus.

② Let  $g_0$  be an Einstein metric, i.e.,

$$\text{Ric}(g_0) = \lambda g_0 \text{ for some } \lambda \in \mathbb{R}.$$

Then  $g(t) = (1-2\lambda t)g_0$  is a sol<sup>n</sup> to  $(RF)$

$$\begin{aligned} \text{as } \partial_t g(t) &= -2\lambda g_0 = -2\text{Ric}(g_0) \\ &= -2\text{Ric}((1-2\lambda t)g_0) \\ &= -2\text{Ric}(g(t)). \end{aligned}$$

$$\text{so } g(t) = 0 \text{ at } t = \frac{1}{2\lambda}.$$

If  $\lambda > 0$  then the solutions are shrinking as  $g(t)$  is shrinking from  $g_0$ .

$\lambda = 0$  static sol<sup>n</sup>

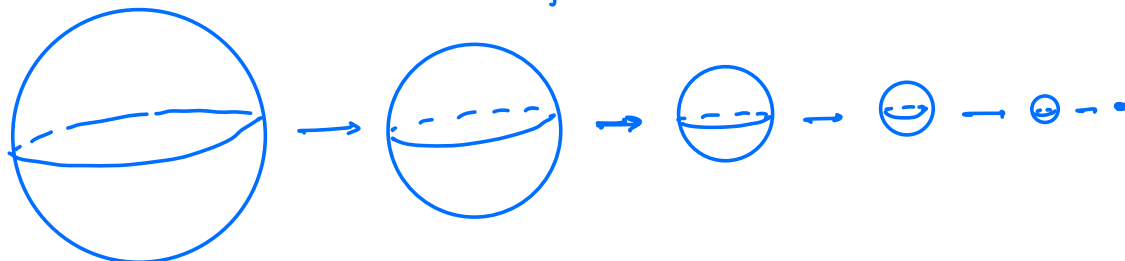
$\lambda < 0$  expanding sol<sup>n</sup>.

A concrete case is that of  $(S^n, g_0)$ . Here

$$\text{Ric}(g_0) = (n-1)g_0 \Rightarrow g(t) = (1-2(n-1)t)g_0$$

is a sol<sup>n</sup> to  $(RF)$ . and the sol<sup>n</sup> exists till  $T = \frac{1}{2(n-1)}$ .

Pictorially, the flow runs by shrinking the sphere until it becomes a point.





Def<sup>n</sup>:-  $g(t)$  is called an **eternal sol<sup>n</sup>** if  $t \in (-\infty, \infty)$ .

$g(t)$  is an **ancient sol<sup>n</sup>** if  $t \in (-\infty, c)$ ,  $c < \infty$ .

(e.g. round sphere)

$g(t)$  is an **immortal sol<sup>n</sup>** if  $t \in (\alpha, \infty)$ ,  $\alpha > -\infty$ .

### Symmetries of RF

\*  $(M, g(t))_{t \in \mathbb{I}}$  is RF  $\rightarrow (M, g_{t-t_0})_{t \in \mathbb{I}+t_0}$  is a RF.

\* Parabolically rescaling a RF gives another RF  
 $\hookrightarrow$  "time scales like (distance)<sup>2</sup>"  
 i.e. if  $g(t)$  is a RF then

$$\hat{g}(x|t) = \lambda g(x, \frac{t}{\lambda}), \quad t \in [0, \lambda T] \text{ is also a RF.}$$

$$\partial_t \hat{g}(x|t) = \lambda \cdot \frac{1}{\lambda} \partial_t g(x, \frac{t}{\lambda}) = -2 \text{Ric}(g(x, \frac{t}{\lambda})) = -2 \text{Ric}(\hat{g}).$$

\* Diffeomorphism invariance. If  $\varphi: M \rightarrow M$  is a diffeo. and  $g(t)$  is a RF then so is  $\varphi^* g(t)$ .

## Ricci flow regarded as a heat eq<sup>n</sup>

Def<sup>n</sup> :- Local coordinates  $(x^i)$  are called harmonic if  $\Delta x^i = 0$ .

$$\therefore 0 = \Delta x^i = g^{jk} (\partial_j \partial_k - \Gamma_{jk}^l \partial_l) x^i = -g^{jk} \Gamma_{jk}^l.$$

Lemma :- For  $p \in M \exists$  harmonic coordinates in some nbd. of  $p$ .

Lemma :- In harmonic coordinates,

$$-2R_{ij} = \underbrace{\Delta(g_{ij})}_{\text{Laplacian of components}} + \underbrace{Q_{ij}(g^{-1}, \partial g)}_{\text{a term involving quadratic expressions in metric inverse } g^{-1} \text{ and } \partial g}.$$

\* Notation :- For tensors  $A, B$

$A * B$  would mean "some linear combination of traces of  $A \otimes B$  w/ coefficients that do not depend on  $A$  or  $B$ ".

So various contractions b/w tensors whose

precise forms are not important.

∴ in harmonic coordinates, indeed

$$\partial_t g_{ij} = \Delta(g_{ij}) + Q_{ij}(g^{-1}, \partial g).$$

Proof of the lemma :-

Recall the formula for  $R_{jk}$  in local coordinates give

$$-2R_{jk} = -2 \left( \partial_g \Gamma_{jk}^p - \partial_j \Gamma_{gk}^p + \Gamma_{jk}^p \Gamma_{gp}^q - \Gamma_{gk}^p \Gamma_{jp}^q \right)$$

$$= -2 \left( \partial_g \left( \frac{1}{2} g^{rs} (\partial_j g_{rk} + \partial_k g_{rj} - \partial_r g_{jk}) \right) \right.$$

$$\left. - \partial_j \left( \frac{1}{2} g^{rs} (\partial_g g_{kr} + \partial_k g_{gs} - \partial_r g_{ks}) \right) \right.$$

$$\left. + \Gamma_{jk}^p \Gamma_{gp}^q - \Gamma_{gk}^p \Gamma_{jp}^q \right)$$

$$= -\partial_g \left( g^{rs} (\partial_j g_{rk} + \partial_k g_{rj} - \partial_r g_{jk}) \right)$$

$$+ \partial_j \left( g^{rs} (\partial_g g_{kr} + \partial_k g_{gs} - \partial_r g_{ks}) \right)$$

$$+ g^{-1} * g^{-1} * \partial g * \partial g$$

(we used the coordinate expression for  $\Gamma_{ij}^k$ )

$$\begin{aligned}
&= g^{lr} \left( -\underline{\partial_q \partial_j g_{lk}} - \partial_q \partial_k g_{lj} + \underline{\partial_q \partial_l g_{jk}} \right) \\
&\quad + g^{lr} \left( \underline{\partial_j \partial_q g_{kl}} + \partial_j \partial_k g_{qr} - \partial_j \partial_l g_{kr} \right) \\
&\quad\quad + g^{-1*} g^{-1*} \partial g^* \partial g \longrightarrow (*)
\end{aligned}$$

— terms cancel

— is the  $\Delta(g_{jk})$  term on contracting  $q$  and  $l$

The remaining 3 terms can be written in terms of partial derivatives of  $\Gamma^s g$ . e.g.

$$g^{lr} \left( -\partial_q \partial_k g_{lj} + \frac{1}{2} \partial_j \partial_k g_{qr} \right)$$

$$= -g^{lr} \partial_k \left( \Gamma_{ql}^s g_{sj} \right)$$

as

$$-g^{lr} \partial_k \left( \Gamma_{ql}^s g_{sj} \right)$$

$$= -g^{lr} \partial_k \left( \frac{1}{2} g^{sl} (\partial_q g_{nl} + \partial_n g_{ql} - \partial_l g_{qr}) \right)$$

$$= -g^{lr} \partial_k \left( \frac{1}{2} (\partial_q g_{lj} + \partial_l g_{qj} - \partial_j g_{ql}) \right)$$

$$= -g^{lr} \left( \frac{1}{2} (\partial_k \partial_q g_{lj} + \partial_k \partial_l g_{qj} - \partial_k \partial_j g_{ql}) \right)$$

$$= -\partial_k \partial_q g_{lj} + \frac{1}{2} \partial_k \partial_j g_{ql}$$

and the remaining term  $g^{rs} (-\partial_j \partial_r g_{ks} + \frac{1}{2} \partial_k \partial_j g_{rs})$

$$\text{in } (*) = -g^{rs} \partial_j (\Gamma_{rs}^s g_{sk})$$

so (\*) becomes

$$-2R_{jk} = \Delta(g_{jk}) - \underbrace{g^{rs} \partial_k (\Gamma_{rs}^s g_{sj}) - g^{rs} \partial_j (\Gamma_{rs}^s g_{sk})}_{+ g^{-1} * g^{-1} * \partial g * \partial g}$$

The underlined \_\_\_\_\_ terms are zero in harmonic coordinates and so

$$-2R_{jk} = \Delta(g_{jk}) + \underbrace{g^{-1} * g^{-1} * \partial g * \partial g}_{Q_{ij}(g^{-1}, \partial g)}$$

□